Learned holographic light transport

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Abstract: Computer-Generated Holography (CGH) algorithms often fall short in matching 7 simulations with results from a physical holographic display. Our work addresses this mismatch 8 by learning the holographic light transport in holographic displays. Using a camera and a 9 holographic display, we capture the image reconstructions of optimized holograms that rely 10 on ideal simulations to generate a dataset. Inspired by the ideal simulations, we learn a 11 complex-valued convolution kernel that can propagate given holograms to captured photographs 12 in our dataset. Our method can dramatically improve simulation accuracy and image quality in 13 holographic displays while paying the way for physically informed learning approaches. 14

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16 1. Introduction

The future of human-computer interactions [1] demands technologies that can display lifelike three-dimensional visuals. An emerging trend, Computer-Generated Holography (CGH),
promises to deliver such realistic visuals in the next-generation displays [2]. However, CGH
algorithms often fall short of achieving high image quality in real life.

The traditional CGH algorithms such as Gerchberg-Saxton method [3] or recently trending approaches such as Stochastic Gradient (SGD) based differentiable methods [4–6] can deliver an outstanding image quality in the simulation environments. However, in an actual holographic display with phase-only modulation, holograms optimized or learned using these ideal holographic light transport models often fail to deliver the same image quality. Identifying causes of mismatch and bridging the gap between the image qualities of simulations and actual experiments are growing scientific research trends in the holography community.

The traditional solutions [7] to address the mismatch aims to find complex residual values that 28 can be added as a regularization term to the ideal holographic light transport [6] or the complex 29 hologram [5]. These techniques to regularize holographic image reconstruction models [5,6] 30 are powerful and effective in practice. In the meantime, researchers have also garnered interest 31 to learn the hologram generation process using deep learning [5,8]. However, their proposed 32 solutions often yield highly complex algorithmic structures and sometimes require a physically 33 demanding experimentation routine. These complex algorithmic structures involve learning 34 components such as Generative Adversarial Networks (GANs) that are not straightforward in 35 tuning and training [6], multi-layer perceptrons that model the nonlinear response of an SLM. 36 which may carry lesser semantic meaning for an optical scientist [5] or characterizing aberrations 37 with Zernike polynomials that requires careful experimentation [5, 9-12]. We ask ourselves if the 38 demand in experimentation load and complex nature of algorithms can be avoided while optical 39 scientists can get more hints towards understanding imperfections in actual holographic displays. 40 With that question in mind, we aim for deriving a new and refined CGH algorithm to improve 41 image quality in actual holographic displays. 42

This work argues that a tailored holographic light transport model for a target holographic display can account for optical aberrations and bridge the gap between simulations and actual holographic displays. We also argue that such a model can avoid intensive experimentation

requirements in display calibration. For this purpose, we propose to learn a single complex-valued 46 point spread function that helps us to propagate input phase-only holograms to the target image 47 plane. Thus, our holographic light transport model convolves an input phase-only hologram 48 with a learned complex-valued point spread function to get to the physically accurate image 49 reconstructions in simulations for a target holographic display. The learning process involves 50 comparing image reconstruction in simulations against experiments using a camera with an 51 actual holographic display. Like any other learning process, we must have a set of data composed 52 of input phase-only holograms and their corresponding image reconstructions in an actual 53 holographic display. We collect such a dataset from our proof-of-concept display prototype using 54 a camera and an ideal holographic light transport based hologram optimization method that is fully 55 differentiable. We show that our learned holographic light transport can dramatically improve 56 simulation accuracy and final image quality in our holographic display. Our key contributions 57 are summarized as follows: 58

Learned holographic light transport. We propose a learned approach for *holographic light transport* to bridge the gap between simulations and experimentation. Our method learns
 a single complex convolutional kernel to reconstruct images in simulation similar to the
 real experiments. Our implementation is fully differentiable. We show that image quality
 results from an actual holographic display can be enhanced with our method while the
 simulations become highly accurate.

Holographic dataset from a proof-of-concept holographic display. In order to be able to train and derive a single complex convolutional kernel, we build a phase-only holographic display. Then, we capture a series of photographs of holographic image reconstructions resulting from holograms optimized using the ideal holographic light transport.

In the following sections, we will first introduce a standard ideal holographic light transport model. Then, we will provide the details of our experimental setup. Finally, we will introduce our technique in learning and provide quantitative results of our method while comparing it against the ideal case.

73 2. Optimizing holograms with ideal holographic light transport

The topic of light transport plays a crucial role in formulating the basis of various domains
including traditional computer graphics [13], architecture [14], biomedical imaging [15], nonline-of-sight imaging [16], three-dimensional printing [17], visible light communications [18],
holographic recording [19], computational displays [20], eye prescription correction [21], eyegaze tracking [22], ophthalmology [23] and many more. Although we cover only display
technologies in this work, an accurate representation method of light transport can potentially
pave the way towards enhancements in many other highlighted applications.

Light transport models used in CGH bases on Rayleigh-Sommerfeld diffraction integrals [24]. This diffraction integral's first solution, the Huygens-Fresnel principle, is expressed as follows:

$$u(x,y) = \frac{1}{j\lambda} \int \int u_0(x,y) \frac{e^{jkr}}{r} \cos(\theta) dx dy,$$
(1)

where resultant field, u(x, y), is calculated by integrating over every point across hologram plane in XY axes, $u_0(x, y)$ represents the optical field in the hologram plane for every point across XY axes, *r* represents the optical path between a selected point in hologram plane and a selected point in target plane, θ represents the angle between these points, *k* represents the wavenumber $(\frac{2\pi}{\lambda})$ and λ represents the wavelength of light. In this model, optical fields, $u_0(x, y)$ and u(x, y), are represented with a complex value,

$$u_0(x, y) = A(x, y)e^{j\phi(x, y)},$$
(2)

where A represents the spatial distribution of amplitude and ϕ represents the spatial distribution of phase across a hologram plane. To simplify our description, we can express the Huygens-Fresnel principle as a superposition of diverging spherical waves originating from a hologram [25]. Perhaps one can also think of the Huygens-Fresnel principle as stamping a complex point-spread function on a target image plane for each point of a hologram while weighting each stamp with its amplitude and phase from its origin.

Calculating Huygens-Fresnel approximation by visiting each point on a hologram one by one
 would consume a large computation and power budget while being slow in processing. Common
 approaches in the literature [26–28] dedicated to near fields (e.g., short distances like 10 cm or
 half a meter) formulates this integral as a convolution operation with a single complex kernel.
 Hence, the common approaches [29] can be expressed as

$$u(x, y) = u_0(x, y) * h(x, y) = \mathcal{F}^{-1}(\mathcal{F}(u_0(x, y))\mathcal{F}(h(x, y))) = U_0(f_x, f_y)H(f_x, f_y),$$
(3)

where h represents a spatially varying complex convolution kernel. The value of the complex
 kernel, h, is typically expressed as

$$h(x, y) = \frac{e^{jkz}}{j\lambda z} e^{\frac{jk}{2z}(x^2 + y^2)},$$
(4)

where z represents the distance between a hologram plane and a target image plane. This ideal model is implemented in a differentiable fashion (refer to odak.learn.wave.classical L81-114) in our fundamental library for optical sciences [30]. The same library hosts differentiable models of various light transport approximations (refer to odak.learn.wave.classical L8-53).

Now that we have established an ideal holographic light transport model as in Equation 3. 106 We can use this holographic light transport model as a forward model that propagates light 107 from a hologram to a target plane. As mentioned earlier, since this model is implemented in 108 code using a modern machine learning library, PyTorch [31], we take advantage of the fact that 109 modern machine learning libraries are capable of automatically differentiating provided functions 110 Differentiation helps to calculate the complex gradient of our forward model's error. In simple 111 terms, for each input phase-only hologram, the resulting image reconstruction can be calculated, 112 and the impact of changing phase values on image reconstruction can be precisely estimated 113 using gradients. This fact helps an optimizer to have meaningful modifications on phase values 114 of a phase-only hologram at the each optimization step. 115

In order to fully realize the described optimization, a loss function is required. For this purpose, we define a loss function, L, using least squared error between a reconstructed image at a target plane, u(x,y), and a target image, t(x,y),

$$L = (u(x, y) - t(x, y))^{2}.$$
 (5)

Note that the loss function described here is the simplest case, and we leave customization of 119 this loss function to meet the application's demands as a future discussion. In addition to a loss 120 function, we would require an optimizer to optimize our phase-only holograms for various targets. 121 We choose to use a Stochastic Gradient Descent based optimization method [32, 33] with a 122 learning rate of 0.1. We ran our optimizer using our ideal forward model for 200 iterations at each 123 hologram calculation. Our hologram optimization method (refer to) is distributed as a part of our 124 fundamental library for optical sciences [30]. We also provide examples at odak.test.learn_sgd 125 for using the optimization method within our library. 126

¹²⁷ Using the described ideal holographic light transport and hologram optimization methodology, ¹²⁸ we calculate phase-only holograms of target images from DIV2K dataset [34]. We resize images in DIV2K dataset to 1920x1080 to match the size of our SLM. We also convert those images to monochrome by taking an average across three color channels. Note that all these images at DIV2K dataset are used only in training (finding the holographic light transport kernel, not in test cases). The calculated image reconstructions perfectly matching the target images in simulations, however they have to be tested against photographs captured from an actual holographic display.
Thus, we will explain how we build a proof-of-concept holographic display in the next section before explaining our final methodology to improve visual quality and realism.

3. Proof-of-concept holographic display

We build a proof-of-concept holographic display to assess image quality of our hologram
optimization methods that uses ideal holographic light transport. We will introduce our learned
holographic light transport in the next section, we will also use the same proof-of-concept
holographic display to assess image quality of our methodology.



Fig. 1. Schematic diagram of our proof-of-concept holographic display prototype used in our experimental setup.

The optical layout of our proof-of-concept holographic display is represented in Figure 1. 141 Following the light from its source, the optical assembly of our proof-of-concept holographic 142 display uses a multi-wavelength laser light source, LASOS MCS4. However, for our experimen-143 tation, we only rely on the working wavelength of 515 nm. A Thorlabs LB1945-A bi-convex 144 lens with a 200 mm focal length lens collimates the output beam of our laser light source. The 145 collimated beam goes through a wire grid linear polarizer, Thorlabs LPVISE100-A, to maintain a 146 polarization aligned with our phase-only Spatial Light Modulator's fast axis (SLM). The linearly 147 polarized collimated beam bounces off an anti-reflection coated Pellicle beamsplitter, Thorlabs 148 BP245B1, towards our 0.90 degrees tilted phase-only SLM, Holoeye Pluto 2.0 (tilted half order). 149 To avoid undiffracted light, we add a horizontal grating to the displayed holograms on our SLM. 150 The horizontally grated hologram, u'_0 can be calculated as 151

$$u_0'(x, y) = \begin{cases} e^{-j(\phi(x, y) + \pi)} & \text{for } x = \text{odd} \\ e^{-j\phi(x, y)} & \text{for } x = \text{even} \end{cases}$$
(6)

where ϕ , the original phase of u_0 , is modified. This way, we steer the location of the reconstructed image in space away from undiffracted light. The tilt angle of our SLM calculated using the 154 diffraction equation formulated as

$$m\lambda = \Delta asin(\theta),\tag{7}$$

where m is the half-order (0.5), Δa is pixel pitch of a SLM and the θ is the angular location of 155 the grated hologram plane. For our system, θ is calculated as 1.80°. So the required tilt angle 156 for the SLM is $\frac{\theta}{2} \approx 0.90^{\circ}$. In the rest of the setup, the phase-modulated beam goes through 157 the Pellicle beamsplitter. In the next stage, the beam passes focusing lenses, a combination 158 of Thorlabs LA1908-A and LB1056-A. A pinhole aperture, Thorlabs SM1D12, follows the 159 lenses at the focal distance of the focusing lenses to avoid undiffracted light. We capture the 160 image reconstructions of our hologram dataset optimized using ideal holographic light transport 161 from our setup with a lensless image sensor, Point Grey GS3-U3-23S6M-C USB 3.0. For each 162 captured image reconstruction, we applied homography correction for the captures, so that we 163 can compare it against a ground-truth image or a simulated reconstruction. The holograms in our 164 work are always reconstructed for a target image plane at 7 cm away from our proof-of-concept 165 holographic display. 166

167 4. Learned holographic light transport

We provide sample photographs showing image reconstructions captured from our proof-ofconcept holographic display in Figure 2. These photographs are a result of holograms optimized using the ideal holographic light transport model. We also provide input holograms and their simulated results for comparison. The visual mismatch between photographs and simulated results provides a good understanding of the image quality issues discussed earlier.



Fig. 2. Mismatch between simulated and experimental results when using ideal holographic light transport. For a given (a) phase-only hologram, A simulated result can provide (b) a perfect image reconstruction, while the same hologram in (c) a real holographic display fail in achieving such image reconstructions as we show in Dataset 1 (Ref. [35]).

To combat this mismatch illustrated in Figure 2, we take advantage of our dataset of photographs from the proof-of-concept prototype and their corresponding optimized holograms that used the

ideal holographic light transport model (Dataset 1 [35]). With a Stochastic Gradient Descent 175 based optimization method [32, 33] and a learning rate of 0.002, we set to learn a complex kernel, 176 $h_1(x, y)$ using the loss function at Equation 5 that will replace the original h(x, y) from the ideal 177 case. This newly optimized $h_l(x, y)$ can be best described as a transfer function that takes an 178 ideal input hologram and provides an image reconstruction similar to the captured photographs 179 in our dataset. The code base of our learning process follows the same optimization described 180 in Section 2 (refer to realistic_holography:optics L87-L137). The phase and amplitude of the 181 learned complex kernel, $h_l(x, y)$, and the ideal complex kernel, h(x, y), are provided in Figure 3 182 for comparison. 183



Fig. 3. A phase and amplitude comparison between complex kernels used in (a) ideal holographic light transport and (b) learned holographic light transport.

184 5. Evaluation

Now that we have a learned transfer function, $h_1(x, y)$, shown in Figure 3, we look into how 185 186 this kernel representing the learned holographic light transport can help us to optimize new holograms. Assume that the learned kernel is more realistic than the original ideal kernel. In that 187 hypothesis, the optimized holograms should lead to image reconstruction results better in terms 188 of image quality in the experimental case. Meantime, we should also expect that the mismatch 189 between simulations and experiment cases to be mitigated. We challenge these assumptions 190 by optimizing holograms using $h_l(x, y)$ instead of h(x, y). In our exploration for optimizing 191 holograms using the learned holographic light transport, we rely on the same process described 192 in Section 2 (refer to realistic_holography:optics L45-85). 193

Image quality. We provided a visual comparison between holograms generated using the ideal holographic light transport and learned holographic light transport in Figure 4. The visual quality of the reconstructed images in our proof-of-concept holographic display using the learned holographic light transport shows a significant improvement over the ideal case. We believe this is because imperfections in our proof-of-concept holographic display are accounted for in our transfer function. We kindly invite the readers to observe the visual difference between the



Fig. 4. A visual comparison between (a) ideal holographic light transport and (b) learned holographic light transport in reconstructing images. Both of the photographs are captured with optimized holograms using corresponding holographic light transport models and our proof-of-concept prototype. Note that target image at both cases are not used in our training set (DIV2K [34]).

ideal transfer function and the learned transfer function provided in Figure 3. Please note the
 asymmetry in the learned kernel, which does not exist in the case of an ideal kernel.

²⁰² The mismatch between simulations and experiments. Our learned holographic light trans-

 $_{203}$ port can approximate a transfer function of our proof-of-concept display accurately (Training L2

loss: 0.0028 and test loss: 0.0034 – learned reconstruction versus captured ground truth images –

note that images are normalized between zero and one). We compare image reconstructions from

²⁰⁶ our simulations with our experimental results from our proof-of-concept holographic display to



Fig. 5. Learned simulation (a) versus real photograph (b). Ideal light transport based hologram optimization estimates unrealistic results in simulation. On the other hand, for a given target image (c), simulations based on learned holographic light transport closely resembles the experimental results.

provide evidence that this is the case. This comparison is sampled in Figure 5. Our simulations' 207 brightness and contrast levels with learned holographic light transport do not truly match our 208 photographs from our experimentation. However, the spatial content in experimental cases 209 resembles the simulated reconstructions closely and even giving us an excellent hint about what 210 to expect in terms of visual quality from a given holographic display. If further tweaking is 21 needed, the brightness mismatch in the ideal and learned cases can be improved by following a 212 manual calibration routine. In the supplementary documentation of work by Choi et al. [36], 213 curious readers can find highly detailed documentation on minimizing the average difference 214 between simulation and a physical prototype by adjusting laser power and exposure time. We 215 have not conducted such a calibration for this work, as we wanted to show the improvement over 216 an uncalibrated system. 217

What did we learn from the learned kernel? The holographic light transport kernel learned 218 within this work (see Figure 3) indicates that the phase and amplitude behaviour of our physical 219 light source is not homogenous in terms of angular emission. The readers may observe this fact by 220 carefully checking the asymmetry of the kernel in Figure 3. The amplitude values are far greater 221 than the ideal kernel, thus suggesting that to get to brighter images, the hologram optimization 222 has to consider this correlation. This fact can be observed in Figure 4 as the dynamic range, 223 and brightness levels are better preserved in the learned method. Note that the learned kernel is 224 the point-spread function of the given holographic display. Thus, resolution characteristics of 225 the holographic display can also be analyzed in the future by studying the limits of a learned 226 point-spread function. Finally, note that a single kernel can only capture a global mean of a 227 general trend in a holographic display. We will discuss how to improve our learned method in 228 the future in the final paragraph of this section. 229

Comparison with the state of the art. The leading state-of-the-art methods [5,6] that bridge 230 the gap between simulations and physical holographic displays consists of convolutional neural 231 networks. Specifically, the work by Peng et al. [5] relies on more than eight million parameters 232 to tune in a training process of neural networks. Many parameters and layers are needed to 233 efficiently realize the correlation between a hologram and a final reconstructed image. Otherwise, 234 the connections between the pixels of a hologram and a final reconstructed image may not be 235 fully identified (locality issue). This locality issue arises from the fact that such models use small 236 kernel sizes. In contrast, our work decreases this number of tunable parameters to half, four 237 million parameters (2x1080x1920 – amplitude and phase), while relying on kernel sizes that is 238

the same as an input image which avoids locality issue. The work by Maimone et al. [37] uses
one-dimensional separable functions for reducing the memory footprint in classical CGH [37]
for ideal complex forms such as quadratic phase functions. Drawing inspiration from that work,
we speculate that further reduction may be possible by storing a parametric form of a learned
complex kernel.

The readers of our work may ask if our approach is the solution that could provide the most 244 remarkable accuracy in bridging the gap between simulations and experiments in holography. 245 We followed a similar approach to the classical model, where a single convolutional kernel 246 formulates the transfer function of light transport. Hence, our approach is accurate as long as 247 a single kernel is reliable to describe the light transport. To improve accuracy further and to 248 have one-to-one matching simulations in the future, we speculate that approaches with spatially 249 varying convolutional kernels can provide more capacity to accommodate for a genuinely realistic 250 simulation. On the other hand, we learn the light transport between a hologram and an image 251 plane. Approaches that provide three-dimensional image reconstructions in CGH require a 252 transfer function representing the relationship between a single hologram and multiple image 253 planes. In our approach, we have to learn the kernel for each plane using a set of images. In the 254 future, a complete form of our approach can potentially be derived where a data set with diverse 255 image reconstruction distances helping to learn a parametric light transport rather than per plane 256 learning. Our work does an excellent job in capturing optical aberrations and imperfections of a 257 holographic display. Our work can be best described as the simplest form of improving realism 258 in CGH algorithms without dealing with complex experimentation or complex algorithmic 259 approaches. 260

261 6. Conclusion

Holographic displays often require tedious effort to optimize holograms for the best possible 262 image quality. We propose a new learned method to address this issue with holographic displays in 263 a simple way. The core of our approach is in a learning procedure that allows one to approximate 264 an accurate holographic light propagation model for a given actual holographic display. With this 265 approach, we can optimize holograms that can dramatically improve image quality concerning a 266 typical ideal holographic light transport model. Our method, in turn, enables a simple yet effective 267 method that does not suffer from the overhead of deriving complex algorithmic approaches while 268 paving the way towards physically informed learning approaches in the holography domain. 269

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274 **Disclosures.** The authors declare no conflicts of interest.

275 Data Availability Statement. The code base discussed in Section 2 and 4 is readily available in the

²⁷⁶ Github:complight/realistic_holography. The generated dataset of this work is also available in the Dataset:

277 Phase-only holograms and photographs.

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